

On the Euler Equations of a Nonhomogeneous Ideal Incompressible Fluid

ノンホモジニアウスな非圧縮性理想流体のオイラー方程式について

伊 藤 成 治*

Shigeharu ITOH

Abstract.

It is shown here that the Cauchy problem for the Euler equations of a nonhomogeneous ideal incompressible fluid has a unique solution for a small time interval. The existence of a solution is established by applying the method of the semi Galerkin approximations.

§1. Introduction.

Consider the system of equations

$$\begin{cases} \rho_t + v \cdot \nabla \rho = 0 \\ \rho [v_t + (v \cdot \nabla) v] + \nabla p = \rho f \\ \operatorname{div} v = 0 \end{cases} \quad (1.1)$$

in $Q_T = \mathbb{R}^3 \times [0, T]$, subject to the initial conditions

$$\begin{cases} \rho|_{t=0} = \rho_0(x) \\ v|_{t=0} = v_0(x). \end{cases} \quad (1.2)$$

Here $f(x, t)$, $\rho_0(x)$ and $v_0(x)$ are given, while the density $\rho(x, t)$, the velocity vector $v(x, t) = (v^1(x, t), v^2(x, t), v^3(x, t))$ and the pressure $p(x, t)$ are unknowns. The system (1.1) describes the motion of a nonhomogeneous ideal incompressible fluid.

Our theorem is the following.

Theorem. *Assume that*

$$\rho_0(x) - \bar{\rho} \in H^3(\mathbb{R}^3) \text{ for some positive constant } \bar{\rho}, \quad (1.3)$$

$$\inf \rho_0(x) \equiv m > 0 \text{ and } \sup \rho_0(x) \equiv M < \infty, \quad (1.4)$$

* 弘前大学教育学部数学科教室

Department of Mathematics, Faculty of Education, Hirosaki University

$$v_0(x) \in H^3(\mathbb{R}^3) \text{ and } \operatorname{div} v_0 = 0 \quad (1.5)$$

and

$$f(x, t) \in L^2(0, T; H^3(\mathbb{R}^3)). \quad (1.6)$$

Then there exists $T^* \in (0, T]$ such that the problem (1.1) and (1.2) has a unique solution (ρ, v, p) which satisfies

$$(\rho - \bar{\rho}, v, \nabla p) \in L^\infty(0, T^*; H^3(\mathbb{R}^3)) \times L^\infty(0, T^*; H^3(\mathbb{R}^3)) \times L^2(0, T^*; H^3(\mathbb{R}^3)). \quad (1.7)$$

§2. Preliminaries.

In this section we shall obtain an a priori estimate for solutions of (1.1) and (1.2). Let (ρ, v, p) be a sufficiently regular solution. We assume, for simplicity, $f=0$. The general case can be treated in the same way.

Lemma 2.1. *If we put $\tilde{\rho} = \rho - \bar{\rho}$, then*

$$\frac{d}{dt} \|\tilde{\rho}(t)\|_3 \leq c_1 \|v(t)\|_3 \|\tilde{\rho}(t)\|_3, \quad (2.1)$$

where c_1 is a positive constant depending only on imbedding theorems

and $\|\cdot\|_K = \|\cdot\|_{H^K(\mathbb{R}^3)}$.

Proof. It follows (1.1)₁ and (1.2)₁ that $\tilde{\rho}$ satisfies the equation

$$\begin{cases} \tilde{\rho}_t + v \cdot \nabla \tilde{\rho} = 0 \\ \tilde{\rho}|_{t=0} = \tilde{\rho}_0(x). \end{cases} \quad (2.2)$$

Applying the operator $D^\alpha = (\partial/\partial x_1)^{\alpha_1} (\partial/\partial x_2)^{\alpha_2} (\partial/\partial x_3)^{\alpha_3}$ to (2.2)₁, multiplying the result by $D^\alpha \tilde{\rho}$, integrating over \mathbb{R}^3 and adding in α with $|\alpha| = (\alpha_1 + \alpha_2 + \alpha_3) \leq 3$, then we have

$$\frac{d}{dt} \|\tilde{\rho}(t)\|_3^2 \leq c_2 \|v(t)\|_3 \|\tilde{\rho}(t)\|_3^2. \quad (2.3)$$

Hence it is easy to see that (2.1) holds.

Q.E.D.

Lemma 2.2. *Put*

$$\Phi(t) = \sum_{j=0}^3 \sum_{|\alpha|=j} \|\sqrt{\rho(t)} D^\alpha v(t)\|_0, \quad (2.4)$$

Then we have

$$\frac{d}{dt} \Phi(t) \leq c_3 [1 + \|\tilde{\rho}(t)\|_3 + \Phi(t)]^5, \quad (2.5)$$

where c_3 is a positive constant depending only on m, M and imbedding theorems.

Proof. We first note that

$$m \leq \rho(x, t) \leq M, \quad (2.6)$$

since we have the representation

$$\rho(x, t) = \rho_0(y(\tau, x, t) |_{\tau=0}), \quad (2.7)$$

where $y(\tau, x, t)$ is the solution of the Cauchy problem

$$\begin{cases} \frac{dy}{d\tau} = v(y, \tau) \\ y |_{\tau=t} = x. \end{cases} \quad (2.8)$$

Therefore we find that

$$\sqrt{m} \|v\|_3 \leq \Phi(t) \leq \sqrt{M} \|v\|_3. \quad (2.9)$$

(i) We multiply (1.1)₂ by v and integrate over \mathbb{R}^3 . Taking account of (1.1)₁ and (1.1)₃, we get

$$\frac{d}{dt} \|\sqrt{\rho} v\|_0 = 0. \quad (2.10)$$

Multiplying by v_t and integrating over \mathbb{R}^3 , then

$$m \|v_t\|_0^2 \leq M \|v\|_1 \|Dv\|_1 \|v_t\|_0, \quad (2.11)$$

where we use the notation $D^k u = \sum_{|\alpha|=k} D^\alpha u$. Thus

$$\|v_t\|_0 \leq c_4 \|v\|_2^2. \quad (2.12)$$

(ii) Apply the operator D^α with $|\alpha|=1$ on each side of (1.1)₂, multiply the result by $D^\alpha v$ and integrate over \mathbb{R}^3 . Noting (2.9) and (2.12), then, similarly to (i), we get

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|\sqrt{\rho} Dv\|_0^2 &\leq \|D\rho\|_2 \|v_t\|_0 \|Dv\|_0 + \|Dv\|_2 \|\sqrt{\rho} Dv\|_0^2 + \|D\rho\|_2 \|v\|_2 \|Dv\|_0^2 \\ &\leq c_5 (\|D\rho\|_2 \|v\|_2^2 \|\sqrt{\rho} Dv\|_0^2 + \|Dv\|_2 \|\sqrt{\rho} Dv\|_0^2 + \|D\rho\|_2 \|v\|_2 \|\sqrt{\rho} Dv\|_0^2). \end{aligned} \quad (2.13)$$

Hence we have

$$\frac{d}{dt} \|\sqrt{\rho} Dv\|_0 \leq c_6 (1 + \|D\rho\|_2 + \|v\|_3)^3. \quad (2.14)$$

If we multiply by $D^\alpha v_t$ and integrate over \mathbb{R}^3 , then we obtain

$$\begin{aligned} m \|Dv_t\|_0^2 &\leq c_7 (\|D\rho\|_2 \|v_t\|_0 \|Dv_t\|_0 + \|v\|_2 \|D^2 v\|_0 \|Dv_t\|_0 \\ &\quad + \|Dv\|_1^2 \|Dv_t\|_0 + \|D\rho\|_2 \|v\|_2 \|Dv\|_0 \|Dv_t\|_0). \end{aligned} \quad (2.15)$$

Therefore

$$\|Dv_t\|_0 \leq c_8 (1 + \|D\rho\|_2) \|v\|_2^2. \quad (2.16)$$

(iii) Making use of the operator D^α with $|\alpha|=2$ in place of the operator D^α with $|\alpha|=1$ and repeating the argument in (ii), we have

$$\frac{d}{dt} \|\sqrt{\rho} D^2 v\|_0 \leq c_9 (1 + \|D\rho\|_2 + \|v\|_3)^4 \quad (2.17)$$

and

$$\|D^2 v_t\|_0 \leq c_{10}(1 + \|D\rho\|_2)^2 \|v\|_3^2. \quad (2.18)$$

(iv) Apply the operator D^α with $|\alpha|=3$ to (1.1)₂, multiply by $D^\alpha v$ and integrate over \mathbb{R}^3 . Then, we get

$$\frac{d}{dt} \|\sqrt{\rho} D^3 v\|_0 \leq c_{11}(1 + \|D\rho\|_2 + \|v\|_3)^5. \quad (2.19)$$

Consequently, it follows from (2.10), (2.14), (2.17) and (2.19) that (2.5) holds.

Q.E.D.

Proposition 2.3. *There exists $T^* \in (0, T]$ such that*

$$\|\tilde{\rho}(t)\|_3 + \|v(t)\|_3 \leq c \quad \text{for } t \leq T^*, \quad (2.20)$$

where c is a positive constant depending only on m , M , $\|\tilde{\rho}_0\|_3$, $\|v_0\|_3$ and imbedding theorems.

Proof. If we set

$$\Psi(t) = 1 + \|\tilde{\rho}(t)\|_3 + \Phi(t), \quad (2.21)$$

then, from Lemma 2.1 and Lemma 2.2, we have a differential inequality

$$\frac{d}{dt} \Psi(t) \leq \tilde{c} \Psi(t)^5, \quad (2.22)$$

where $\tilde{c} = c_1 + c_3$. Thus we conclude that

$$\Psi(t) \leq \Psi(0) (1 - 4\tilde{c}t\Psi(0)^4)^{-1/4} \quad \text{as long as } t < (4\tilde{c}\Psi(0)^4)^{-1}. \quad (2.23)$$

Q.E.D.

§3. Proof of Theorem.

We solve the problem (1.1) and (1.2) by applying the semi Galerkin method with the basis $\{\varphi_k(x)\}$ in $H^4(\mathbb{R}^3) \cap J$, where $J = \{u \in \{C_0^\infty(\mathbb{R}^3)\}^3 : \operatorname{div} u = 0\}$. Let us look for $\rho_n(x, t)$ and

$$v_n(x, t) = \sum_{j=1}^n a_{nj}(t) \varphi_j(x) \quad (3.1)$$

satisfying

$$\begin{cases} \rho_{nt} + v_n \cdot \nabla \rho_n = 0 \\ ((\rho_n [v_{nt} + (v_n \cdot \nabla) v_n], \varphi_k)) = 0, \quad k=1, \dots, m \\ \rho_n|_{t=0} = \rho_0(x) \\ v_n|_{t=0} = \sum_{j=1}^n a_j \varphi_j(x), \quad a_j = \int_{\mathbb{R}^3} v_0 \cdot \varphi_j dx, \end{cases} \quad (3.2)$$

where $((\cdot, \cdot))$ stands for the scalar product in $H^3(\mathbb{R}^3)$.

If we multiply (3.2)₂ by $a_{nk}(t)$ and add in k , then we obtain

$$((\rho_n [v_{nt} + (v_n \cdot \nabla) v_n], v_n)) = 0. \quad (3.3)$$

Therefore, similar to §2, a priori estimates

$$m \leq \rho_n(x, t) \leq M \quad (3.4)$$

and

$$\|\tilde{\rho}_n(t)\|_3 + \|v_n(t)\|_3 \leq c \quad \text{for } t \leq T^* \quad (3.5)$$

hold. These estimates guarantee the unique solvability of the problem (3.2), and, furthermore, permit to pass to the limit using the standard compactness arguments (cf. [1], [2], [3]).

Hence we can verify the existence of a unique solution of the problem (1.1) and (1.2) as well as the applicability of the inequalities (2.6) and (2.20).

This completes the proof.

References

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(1993.5.12受理)