ON THE VANISHING VISCOSITY IN THE CAUCHY PROBLEM FOR EQUATIONS OF A NONHOMOGENEOUS INCOMPRESSIBLE FLlTID II

ノンホモジニアウスな非圧縮性流体の方程式に対する コーシー問題における粘性消滅についてII

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ABSTRACT We investigate the Cauchy problem for Euler and Navier-Stokes equations of a nonhomogeneous incompressible fluid in \mathbb{R}^{3} . The unique solvability on a small time interval independent of viscosity is proved, and moreover, it is shown that the solution of Navier-Stokes equations converges in some Hilbert space to the one of Euler equations as viscosity tends to zero.

Key words : Incompressible fluid, Navier-Stokes equations, Euler equations, Vanishing viscosity

1. INTRODUCTION

 \mathcal{L}

We consider the system of equations

(1.1)
$$
\begin{cases} \rho_t + v \cdot \nabla \rho = 0, \\ \rho[v_t + (v \cdot \nabla) v] + \nabla p = \mu \Delta v + \rho f, \\ \text{div} \, v = 0, \end{cases}
$$

in $Q_T = \mathbb{R}^3 \times [0, T]$, $T > 0$, subject to the initial conditions

(1.2)
$$
\begin{cases} \rho|_{t=0} = \rho_0(x), \\ v|_{t=0} = v_0(x). \end{cases}
$$

Here $f(x, t)$ is a given vector field of external forces, while the density $\rho(x, t)$, the velocity vector $v(x,t)$ and the pressure $p(x,t)$ are the unknowns. The viscosity coefficient μ is

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assumed to be a nonnegative constant.

In these equations, $p(x,t)$ is automatically determined (up to a function of t) by $\rho(x,t)$ and $v(x,t)$, namely, by solving the equation

(1.3)
$$
\operatorname{div}(\rho^{-1}\nabla p) = \operatorname{div}(f - (v \cdot \nabla)v + \mu \rho^{-1} \Delta v).
$$

Thus we mention (ρ, v) when we talk about the solution of problem (1.1) , (1.2) .

Compared with the previous paper $[2]$, in which the similar results were proved, we discuss the problem under the weaker assumptions to given data.

The purpose of this paper is to prove

Theorem. Let $0 \leq \mu \leq 1$, and assume that

(1.4)
$$
\rho_0(x) \in C^0(\mathbb{R}^3), \ \nabla \rho_0(x) \in H^2(\mathbb{R}^3), \ 0 < m \leq \rho_0(x) \leq M < \infty,
$$

(1.5)
$$
v_0(x) \in H^3(\mathbb{R}^3), \text{ div } v_0 = 0,
$$

(1.6)
$$
f(x,t) \in L^2(0,T;H^3(\mathbb{R}^3)).
$$

Then there exists $T_0 \in (0, T]$ *independent of* μ *such that problem* (1.1), (1.2) *has a unique solution* $(\rho, v)(x, t)$ *which satisfies*

 (1.7)

$$
\rho(x,t) \in C^{0}(\mathbb{R}^{3} \times [0, T_{0}]), \ \nabla \rho(x,t) \in C^{0}([0, T_{0}]; H^{2}(\mathbb{R}^{3})) \ , 0 \leq m \leq \rho(x,t) \leq M < \infty,
$$

(1.8)
$$
v(x,t) \in C^0([0,T_0]; H^3(\mathbb{R}^3)).
$$

Furthrmore, let(ρ^0 , v^0) *be the solution of problem* (1.1), (1.2) *with* $\mu = 0$ *and* (ρ^{μ} , v^{μ}) *the one* with $\mu > 0$, then we have

(1.9)
$$
\sup_{0 \leq t \leq T_0} [\| (\rho^0 - \rho^{\mu})(t) \|_2 + \| (v^0 - v^{\mu})(t) \|_2] \to 0 \text{ as } \mu \to 0,
$$

where $\|\cdot\|_{k} = \|\cdot\|_{H^{k}(\mathbb{R}^{3})}$ *.*

2. PRELIMINARIES

In this section we establish several *a priori* estimates for solutions of problem (1.1),

(1.2). Let $(\rho, v)(x, t)$ be a sufficiently regular solution. Hereafter C stands for the generic constant independent of μ .

Lemma 2.1. *Let*

(2.1)
$$
\Psi(t) = \int_0^t [1 + ||\nabla \rho(s)||_2^2 + ||v(s)||_3^2]^2 ds,
$$

then the estimates

$$
(2.2) \t\t\t\t\t m \leq \rho \left(x, t\right) \leq M
$$

and

(2.3)
$$
\|\nabla \rho(t)\|_2^2 \le \|\nabla \rho_0\|_2^2 + C \Psi(t)
$$

hold.

Proof It is well-known that, according to the classical method of characteristics, the solution of problem (1.1) ₁, (1.2) ₁ is given by

$$
\rho(x,t) = \rho_0(y(\tau,x,t))_{\tau=0},
$$

where $y(\tau, x, t)$ is the solution of the Cauchy problem

(2.5)
$$
\begin{cases} \frac{dy}{d\tau} = v(y, \tau), \\ y|_{\tau = t} = x. \end{cases}
$$

From this, the estimate (2.2) results.

Next let us establish (2.3). Apply the operator D_x^{α} on each side of (1.1)₁. Multiplying the result by $D_x^{\alpha} \rho$, integrating over \mathbb{R}^3 and summing over $|\alpha| = 1, 2, 3$, we have the equality

$$
(2.6) \qquad \frac{1}{2} \frac{d}{dt} \|\nabla \rho(t)\|_2^2 = -\sum_{|\alpha|=1}^3 \int_{\mathbb{R}^3} (v \cdot \nabla D_x^{\alpha} \rho) D_x^{\alpha} \rho dx
$$

$$
-\sum_{|\alpha|=1}^3 \sum_{0 \leq \beta \leq \alpha} {\alpha \choose \beta} \int_{\mathbb{R}^3} (D_x^{\beta} v \cdot \nabla D_x^{\alpha-\beta} \rho) D_x^{\alpha} \rho dx.
$$

The first term of the right hand side is zero, by integration by parts, since div $v = 0$. The second term can be estimated as follows:

$$
\sum_{|\alpha|=1}^3\sum_{0<\beta\leq\alpha}\binom{\alpha}{\beta}\Big|\int_{\mathbb{R}^3}(D_x^{\beta}v\cdot\nabla D_x^{\alpha-\beta}\rho)D_x^{\alpha}\rho dx\Big|\leq C\|\nabla v(t)\|_2\|\nabla\rho(t)\|_2^2.
$$

Hence we get

(2.7)
$$
\frac{d}{dt} \|\nabla \rho(t)\|_2^2 \leq C \|\nabla v(t)\|_2 \|\nabla \rho(t)\|_2^2,
$$

and thus (2.3) is obtained.

Lemma 2.2. *If we put*

(2.8)
$$
A = 1 + \|\nabla \rho_0\|_2^2 + \|v_0\|_3^2 + \int_0^T \|f(t)\|_3^2 dt,
$$

then we have the inequality

(2.9)
$$
\|v(t)\|_3^2 + \int_0^t \|v_t(s)\|_2^2 ds + \mu \int_0^t \|\nabla v(s)\|_3^2 ds \leq C \left[A + \Psi(t)\right]^3.
$$

Proof. By applying the operator D_x^{α} on both sides of $(1.1)_2$, we obtain the equation

$$
(2.10) \qquad \rho \left[D_x^{\alpha} v_t + (v \cdot \nabla) D_x^{\alpha} v \right] + \nabla D_x^{\alpha} \rho
$$

$$
= \mu \triangle D_x^{\alpha} v - \sum_{0 < \beta \leq \alpha} {\alpha \choose \beta} D_x^{\beta} \rho D_x^{\alpha - \beta} v_t
$$

$$
- \sum_{0 < \beta \leq \alpha} {\alpha \choose \beta} \sum_{0 \leq \gamma \leq \beta} {\beta \choose \gamma} D_x^{\gamma} \rho (D_x^{\beta - \gamma} v \cdot \nabla) D_x^{\alpha - \beta} v
$$

$$
+ \sum_{0 \leq \beta \leq \alpha} {\alpha \choose \beta} D_x^{\beta} \rho D_x^{\alpha - \beta} f.
$$

Step 1. Multiplying (2.10) by $D_x^{\alpha} v$ and integrating over \mathbb{R}^3 , then by making use of

 $(1.1)_{1,3}$, (2.2) , we have the inequalities

(2.11;0)
$$
\frac{1}{2} \frac{d}{dt} || \sqrt{\rho} v ||_0^2 + \mu || \nabla v ||_0^2 \leq C ||f||_0 ||v||_0
$$

and for *k=1,2,3,*

(2.11;k)

$$
\begin{aligned}\n&\frac{1}{2} \frac{d}{dt} \|\sqrt{\rho} \nabla^k v\|_0^2 + \mu \|\nabla^{k+1} v\|_0^2 \\
&\leq C \left[\|\nabla \rho\|_2 \|v_t\|_{k-1} \|v\|_3 + (1 + \|\nabla \rho\|_2) (\|f\|_3 + \|v\|_3^2) \|v\|_3 \right] \\
&\leq C \left[\|\nabla \rho\|_2^2 \|v\|_3^2 + (1 + \|\nabla \rho\|_2) (\|f\|_3 + \|v\|_3^2) \|v\|_3 \right] + \frac{m}{2} \|v_t\|_{k-1}^2.\n\end{aligned}
$$

Step 2. If we multiply (2.10) by $D_x^{\alpha} v_t$ and integrate over \mathbb{R}^3 , then we obtain

(2.12;0)
$$
m \|v_t\|_0^2 + \frac{\mu}{2} \frac{d}{dt} \|\nabla v\|_0^2 \leq C \left[\|v\|_2^4 + \|f\|_0^2 \right] + \frac{m}{2} \|v_t\|_0^2
$$

and for $k=1, 2$,

(2.12;k)
$$
m \|\nabla^k v_t\|_0^2 + \frac{\mu}{2} \frac{d}{dt} \|\nabla^{k+1} v\|_0^2
$$

$$
\leq C \left[\|\nabla_{\boldsymbol{\rho}} \|_{2}^{2} \| v_{t} \|_{k-1}^{2} + (1 + \|\nabla_{\boldsymbol{\rho}} \|_{2})^{2} (\|f\|_{2}^{2} + \|v\|_{3}^{2}) \right] + \frac{m}{2} \|\nabla^{k} v_{t} \|_{0}^{2},
$$

which mean

(2.13;0)
$$
m \|v_t\|_0^2 + \mu \frac{d}{dt} \|\nabla v\|_0^2 \leq C \left[\|v\|_2^4 + \|f\|_0^2 \right]
$$

and for $k=1, 2$,

$$
(2.13;k) \t\t m \|\nabla^k v_t\|_0^2 + \mu \frac{d}{dt} \|\nabla^{k+1} v\|_0^2
$$

$$
\leq C \left[\|\nabla_{\boldsymbol{\rho}} \|_{2}^{2} \| v_{t} \|_{k-1}^{2} + (1 + \|\nabla_{\boldsymbol{\rho}} \|_{2})^{2} (\|f\|_{2}^{2} + \|v\|_{3}^{2}) \right].
$$

Step 3. Adding (2.11;1) to (2.13;0), we get

(2.14)
$$
\frac{d}{dt} \|\sqrt{\rho} \nabla v\|_0^2 + 2\mu \|\nabla^2 v\|_0^2 + m \|v_t\|_0^2 + 2\mu \frac{d}{dt} \|\nabla v\|_0^2
$$

$$
\leq C \left[(1 + \|\nabla \rho\|_2^2 + \|v\|_3^2)^2 + \|f\|_3^2 \right],
$$

and noting that $0 \le \mu \le 1$,

$$
(2.15) \t\t\t\t\t\|\nabla v\|_0^2 + \int_0^t \|v_t\|_0^2 ds + \mu \int_0^t \|\nabla^2 v\|_0^2 ds \leq C\left[A + \Psi(t)\right].
$$

Step 4. If we add (2.11;2) to (2.13;1), then we obtain

(2.16)
$$
\frac{d}{dt} \|\sqrt{\rho} \nabla^2 v\|_0^2 + 2\mu \|\nabla^3 v\|_0^2 + m \|\nabla v_t\|_0^2 + 2\mu \frac{d}{dt} \|\nabla^2 v\|_0^2
$$

$$
\leq C \left[(1 + \|\nabla \rho\|_2^2 + \|v\|_3^2)^2 + (1 + \|\nabla \rho\|_2^2) (\|v_t\|_0^2 + \|f\|_3^2) \right].
$$

Hence, due to (2.3) and (2.15),

Step 5. Add (2.11;3) to (2.13;2). Similarly to the above, we have

(2.18)
$$
\frac{d}{dt} \|\sqrt{\rho} \nabla^3 v\|_0^2 + 2\mu \|\nabla^4 v\|_0^2 + m \|\nabla^2 v_t\|_0^2 + 2\mu \frac{d}{dt} \|\nabla^3 v\|_0^2
$$

$$
\leq C \left[(1 + \|\nabla_{\boldsymbol{\rho}}\|_{2}^{2} + \|v\|_{3}^{2})^{2} + (1 + \|\nabla_{\boldsymbol{\rho}}\|_{2}^{2}) \left(\|v_{t}\|_{1}^{2} + \|f\|_{3}^{2} \right) \right],
$$

and thus,

Consequently, from $(2.11;0)$, (2.15) , (2.17) , (2.19) , we find that (2.9) is deduced.

Lemma 2.3. *There exists* $T_0 \in (0, T]$ *independent of* μ *such that*

$$
\Psi(t) \le A \text{ for } t \le T_0.
$$

Proof. If we set $Y(t) = \Psi(t) + A$, then from Lemma 2.1, 2.2, we have a differential inequality

$$
(2.21) \t\t \t\t \frac{d}{dt} Y(t) \le C Y^3(t).
$$

Therefore we find that

(2.22)
$$
Y(t) \le \frac{A}{\sqrt{1 - 2CA^2 t}} \text{ provided } t < (2CA^2)^{-1},
$$

and thus

$$
\Psi(t) \le A \text{ for } t \le T_0 = \frac{3}{8CA^2}.
$$

3. PROOF OF THEOREM

First, we note that from Lemma 2.1, 2.2, 2.3, the estimate

(3.1)

$$
\sup_{0\leq t\leq T_0} \left[\|\nabla \rho(t)\|_2^2 + \|v(t)\|_3^2 \right] + \int_0^{T_0} \|v_t(t)\|_2^2 dt + \mu \int_0^{T_0} \|\nabla v(t)\|_3^2 dt \leq C
$$

is valid.

The unique solvability of problem (1.1), (1.2) is proved with semi Galerkin method based on (3.1) . As this process is parallel with that of $[1, Chapter 3]$, we omit it here and restrict ourselves to establish (1.9).

Subtracting (1.1) with $\mu > 0$ from (1.1) with $\mu = 0$, we get the following linear system of equations for $\sigma = \rho^0 - \rho^\mu$, $w = v^0 - v^\mu$ and $q = p^0 - p^\mu$:

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(3.2)
\n
$$
\begin{cases}\n\sigma_t + v^{\mu} \cdot \nabla \sigma = -w \cdot \nabla \rho^0, \\
\rho^{\mu} [w_t + (v^{\mu} \cdot \nabla) w] + \nabla q = -\rho^{\mu} (w \cdot \nabla) v^0 + (\nabla \rho^0 / \rho^0) \sigma - \mu \Delta v^{\mu}, \\
\text{div } w = 0, \\
\sigma|_{t=0} = 0, \\
w|_{t=0} = 0.\n\end{cases}
$$

In the same way for getting *a priori* estimates, we have from (3.1),

(3.3)
$$
\|\sigma(t)\|_2^2 \leq C \int_0^t \|w(s)\|_2^2 ds
$$

and

(3.4)
$$
\|w(t)\|_2^2 \leq C \int_0^t (\|\sigma(s)\|_2^2 + \|w(s)\|_2^2) ds + \mu CT_0.
$$

Hence, by Gronwall's inequality, we find that

(3.5)
$$
\|\sigma(t)\|_{2}^{2} + \|w(t)\|_{2}^{2} \leq \mu CT_{0} \exp(CT_{0}).
$$

This completes the proof of Theorem.

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